Traveling waves of the spread of avian influenza

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TRAVELING WAVES OF THE SPREAD OF AVIAN INFLUENZA

ZHI-CHENG WANG, JIANHONG WU, AND RONGSONG LIU

(Communicated by Yingfei Yi)

ABSTRACT. This paper gives a proof for the existence and nonexistence of traveling wave solutions of a reaction-convection epidemic model for the spatial spread of H5N1 avian influenza involving a wide range of bird species and environmental contamination. The threshold condition for the existence of traveling waves coincides with the basic reproduction number exceeding one. The existence of wave solutions is obtained by constructing an invariant cone of initial functions defined on a large spatial domain, applying a fixed point theorem on this cone and then a limiting argument. The invariant cone is based on the information of initial growth pattern of the epidemic and the final size estimation during the entire course of the outbreak.

1. INTRODUCTION

We consider the system of reaction diffusion equations

(1.1)
$$
\begin{cases}
\dfrac{\partial S_c}{\partial t} = -\beta_c S_c V - \alpha_c \dfrac{S_c I_c}{N_c} - v_c \dfrac{\partial S_c}{\partial x} + D_c \dfrac{\partial^2 S_c}{\partial x^2}, \\[2mm]
\dfrac{\partial S_w}{\partial t} = -\beta_w S_w V - \alpha_{ew} \dfrac{S_w E_w}{N_w} - \alpha_{iw} \dfrac{S_w I_w}{N_w} - v_w \dfrac{\partial S_w}{\partial x} + D_w \dfrac{\partial^2 S_w}{\partial x^2}, \\[2mm]
\dfrac{\partial I_c}{\partial t} = \beta_c S_c V + \alpha_c \dfrac{S_c I_c}{N_c} - d_{ic} I_c - v_c \dfrac{\partial S_c}{\partial x}, \\[2mm]
\dfrac{\partial E_w}{\partial t} = \beta_w S_w V + \alpha_{ew} \dfrac{S_w E_w}{N_w} + \alpha_{iw} \dfrac{S_w I_w}{N_w} - \mu_w E_w - v_w \dfrac{\partial E_w}{\partial x} + D_w \dfrac{\partial^2 E_w}{\partial x^2}, \\[2mm]
\dfrac{\partial I_w}{\partial t} = \mu_w E_w - d_{iw} I_w, \\[2mm]
\dfrac{\partial I_d}{\partial t} = \beta_d \left(N_d - I_d\right) V + \alpha_d \left(N_d - I_d\right) I_d / N_d - \gamma_d I_d - v_d \dfrac{\partial I_d}{\partial x} + D_d \dfrac{\partial^2 I_d}{\partial x^2}, \\[2mm]
\dfrac{\partial V}{\partial t} = r_c I_c + r_{ew} E_w + r_{iw} I_w + r_d I_d - \left(d_v + d_n\right) V - v_v \dfrac{\partial V}{\partial x} + D_v \dfrac{\partial^2 V}{\partial x^2},
\end{cases}
$$

which was introduced in [9] to describe the spatio-temporal spread of H5N1 avian influenza in an ecosystem involving a wide range of bird species: poultry (c), wild birds (w) which are susceptible to and die after H5N1 infection, and wild birds (d) which are susceptible to but can survive after H5N1 infection. Birds are further

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stratified by their disease status as susceptible (S_c, S_w and S_d), expected (E_w), and infected (I_c, I_w and I_d). Note that the w class wild birds can fly some distances even after exposure to the virus, and hence we have included the expected class E_w. Note also that the total number, N_d, of d class birds is assumed to be a constant. The model also involves the virus (v) in the environment, and exposed/infected birds may contribute to environmental contamination. In the model, D_j and v_j, with $j = w, d, c, v$, the diffusion and convection coefficients of the category j bird/virus, mass action is used to describe the transmission from the virus in the environment to birds, while the standard incidence is used to describe the transmission between birds ($N_c = S_c + I_c$, $N_w = S_w + E_w + I_w$). Finally, the virus production is proportional to the number of infected birds.

The basic reproduction number of the corresponding ODE system is given by $\mathcal{R}_0 = \rho\left(\mathcal{F}\mathcal{V}^{-1}\right)$, where ρ is the spectral radius of a matrix. The matrices $\mathcal{F}$ and $\mathcal{V}$ are relevant to the linearization of the corresponding ODE system of (1.1) at the disease free state $E_0\left(S_c, S_w, I_c, E_w, I_w, I_d, V\right) = (S_{c0}, S_{w0}, 0, 0, 0, 0, 0)$ and are given by

$$\mathcal{F} = \begin{pmatrix} \alpha_c & 0 & 0 & 0 & \beta_c S_{c0} \\ 0 & \alpha_{ew} & \alpha_{iw} & 0 & \beta_w S_{w0} \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \alpha_d & \beta_d N_d \\ r_c & r_{ew} & r_{iw} & r_d & 0 \end{pmatrix}, \mathcal{V} = \begin{pmatrix} d_{ic} & 0 & 0 & 0 & 0 \\ 0 & \mu_w & 0 & 0 & 0 \\ 0 & -\mu_w & d_{iw} & 0 & 0 \\ 0 & 0 & 0 & \gamma_d & 0 \\ 0 & 0 & 0 & 0 & d_v + d_n \end{pmatrix}.$$

The matrices $\mathcal{F}$ and $\mathcal{V}$ and the reproduction number determine the initial growth pattern of the corresponding ODE system. A preliminary analysis of the model was conducted in [9], where the existence of traveling waves was formally studied and the linkage between the minimal wave speed and the disease propagation rate as well as its implication for the effectiveness of different intervention strategies was described numerically. Here, we provide a rigorous proof for the existence or nonexistence of nontrivial traveling wave solutions depending on the size of the basic reproduction number $\mathcal{R}_0$. In addition, we show that when $\mathcal{R}_0 > 1$, there exists an s^* such that (1.1) admits a nontrivial traveling wave solution for every wave speed $s > s^*$. We refer to [9] for more detailed discussions of the biological relevance of these results.

The basic idea to prove the existence of nontrivial traveling wave solutions is to first construct an appropriately invariant cone of initial functions defined in a large but bounded domain, then apply a fixed point theorem on this cone for the relevant solution operators, and finally to pass to the unbounded spatial domain $\mathbb{R}$ by a limiting argument. This method is motivated by [4, 5], where the existence and nonexistence of traveling wave solutions for some infection-age structured epidemic models with diffusion are studied. Adaptation of the approach developed in [4, 5] to our model is highly nontrivial, as the multiple bird species and virus involved make the construction of an invariant cone very difficult. Here we successfully construct such a cone by using the initial growth pattern and the final size outbreak which are encoded by the matrices $\mathcal{F}$ and $\mathcal{V}$, the reproduction number, and the spatial diffusion. Relevant to our work here are [8, 10], and we hope our work provides a framework for more general diffusive epidemic models with species diversity.

2. Nonexistence of traveling wave solutions

A traveling wave of (1.1) is a solution with the form $(S_c(x+st), I_c(x+st),$ $S_w(x+st), E_w(x+st), I_w(x+st), I_d(x+st), V(x+st))$. So, with the wave variable $\xi = x + st$, we have

$$(2.1) \quad \begin{cases} (s+v_c)\, S'_c = -\beta_c S_c V - \alpha_c S_c I_c /\, N_c + D_c S_c, \\ (s+v_w)\, S'_w = -\beta_w S_w V - \alpha_{ew} S_w E_w /\, N_w - \alpha_{iw} S_w I_w /\, N_w + D_w S''_w, \\ (s+v_c)\, I'_c = \beta_c S_c V + \alpha_c S_c I_c /\, N_c - d_{ic} I_c, \\ (s+v_w)\, E'_w = \beta_w S_w V + S_w\, [\alpha_{ew} E_w + \alpha_{iw} I_w] /\, N_w - \mu_w E_w + D_w E''_w, \\ s I'_w = \mu_w E_w - d_{iw} I_w, \\ (s+v_d)\, I'_d(\xi) = \beta_d\, (N_d - I_d)\, V + \alpha_d\, (N_d - I_d)\, I_d /\, N_d - \gamma_d I_d + D_d I''_d, \\ (s+v_v)\, V' = r_c I_c + r_{ew} E_w + r_{iw} I_w + r_d I_d - (d_v + d_n)\, V + D_v V''. \end{cases}$$

Theorem 2.1. *Assume that $\mathcal{R}_0 = \rho\left(\mathcal{F}\mathcal{V}^{-1}\right) < 1$. Then for any $c > 0$, the trivial solution $(S_c \equiv S_{c0}, S_w \equiv S_{w0}, I_c \equiv 0, E_w \equiv 0, I_w \equiv 0, I_d \equiv 0, V \equiv 0)$ is the unique nonnegative and bounded solution of (2.1) satisfying*

$$(2.2) \qquad\qquad S_c\,(-\infty) = S_{c0}, \;\; S_w\,(-\infty) = S_{w0},$$

$$(2.3) \qquad I_c\,(-\infty) = E_w\,(-\infty) = I_w\,(-\infty) = I_d\,(-\infty) = V\,(-\infty) = 0.$$

Proof. Note that if $(I_c, E_w, I_w, I_d, V) \equiv \mathbf{0}$, then $S_c \equiv S_{c0}$ and $S_w \equiv S_{w0}$. Assume that (2.2) and (2.3) hold and (I_c, E_w, I_w, I_d, V) is not identically zero. An application of the fluctuation lemma ([7]) yields that $S'_c\,(-\infty) = S'_w\,(-\infty) = 0$. Consequently, we can show that $S'_c\,(x) \le 0$ and $S'_w\,(x) \le 0$ for $x \in \mathbb{R}$. Thus, we have $S_c\,(x) \le S_{c0}$ and $S_w\,(x) \le S_{w0}$ for $x \in \mathbb{R}$. In particular, we have $I_d(x) \le N_d$ for $x \in \mathbb{R}$. Note that the system of equations for (I_c, E_w, I_w, I_d, V) is equivalent to

$$\begin{cases} I_c\,(\xi) = \int_{-\infty}^{\xi} \frac{d_{ic}}{s+v_c} e^{-\frac{d_{ic}}{s+v_c}(\xi-t)} \frac{1}{d_{ic}}\, [\beta_c S_c\,(t)\, V\,(t) + \alpha_c S_c\,(t)\, I_c\,(t)/\, N_c\,(t)]\, dt, \\ E_w\,(\xi) = \int_{-\infty}^{\xi} \frac{\mu_w}{\rho_w} e^{\lambda_w^-(\xi-t)} \frac{1}{\mu_w} H_w\,(t)\, dt + \int_{\xi}^{\infty} \frac{\mu_w}{\rho_w} e^{\lambda_w^+(\xi-t)} \frac{1}{\mu_w} H_w\,(t)\, dt, \\ I_w\,(\xi) = \int_{-\infty}^{\xi} \frac{d_{iw}}{s} e^{-\frac{d_{iw}}{s}(\xi-t)} \frac{\mu_w}{d_{iw}} E_w(t)\, dt, \\ I_d\,(\xi) = \int_{-\infty}^{\xi} \frac{\gamma_d}{\rho_d} e^{\lambda_d^-(\xi-t)} \frac{1}{\gamma_d} H_d\,(t)\, dt + \int_{\xi}^{\infty} \frac{\gamma_d}{\rho_d} e^{\lambda_d^+(\xi-t)} \frac{1}{\gamma_d} H_d\,(t)\, dt, \\ V\,(\xi) = \int_{-\infty}^{\xi} \frac{d_v+d_n}{\rho_v} e^{\lambda_v^-(\xi-t)} \frac{1}{d_v+d_n} H_v\,(t)\, dt + \int_{\xi}^{\infty} \frac{d_v+d_n}{\rho_v} e^{\lambda_v^+(\xi-t)} \frac{1}{d_v+d_n} H_v\,(t)\, dt, \end{cases}$$

where

$$\lambda_w^{\pm} = \frac{(s+v_w) \pm \sqrt{(s+v_w)^2 + 4D_w\mu_w}}{2D_w}, \quad \lambda_d^{\pm} = \frac{(s+v_d) \pm \sqrt{(s+v_d)^2 + 4D_d\gamma_d}}{2D_d},$$

$$\lambda_v^{\pm} = \frac{(s+v_v) \pm \sqrt{(s+v_v)^2 + 4D_v\,(d_v+d_n)}}{2D_v},$$

$$\rho_w = \lambda_w^+ - \lambda_w^-, \quad \rho_d = \lambda_d^+ - \lambda_d^-, \quad \rho_v = \lambda_v^+ - \lambda_v^-,$$

$$H_w\,(t) := \beta_w S_w\,(t)\, V\,(t) + \alpha_{ew} S_w\,(t)\, E_w\,(t)/\, N_w\,(t) + \alpha_{iw} S_w\,(t)\, I_w\,(t)/\, N_w\,(t),$$

$$H_d\,(t) := \beta_d\,(N_d - I_d\,(t))\, V\,(t) + \alpha_d\,(N_d - I_d\,(t))\, I_d\,(t)/\, N_d,$$

$$H_v\,(t) := r_c I_c\,(t) + r_{ew} E_w\,(t) + r_{iw} I_w\,(t) + r_d I_d\,(t).$$

Therefore,

$$
\begin{cases}
I_c\left(\xi\right) \leq \int_{-\infty}^{\xi} \frac{d_{ic}}{s+v_c} e^{-\frac{d_{ic}}{s+v_c}(\xi-t)} \frac{1}{d_{ic}} \left[\beta_c S_{c0} V\left(t\right) + \alpha_c I_c\left(t\right)\right] dt, \\
E_w\left(\xi\right) \leq \int_{-\infty}^{\xi} \frac{\mu_w}{\rho_w} e^{\lambda_w^-(\xi-t)} \frac{1}{\mu_w} \left[\beta_w S_{w0} V\left(t\right) + \alpha_{ew} E_w\left(t\right) + \alpha_{iw} I_w\left(t\right)\right] dt \\
\qquad + \int_{\xi}^{\infty} \frac{\mu_w}{\rho_w} e^{\lambda_w^+(\xi-t)} \frac{1}{\mu_w} \left[\beta_w S_{w0} V\left(t\right) + \alpha_{ew} E_w\left(t\right) + \alpha_{iw} I_w\left(t\right)\right] dt, \\
I_w\left(\xi\right) = \int_{-\infty}^{\xi} \frac{d_{iw}}{s} e^{-\frac{d_{iw}}{s}(\xi-t)} \frac{\mu_w}{d_{iw}} E_w dt, \\
I_d\left(\xi\right) \leq \int_{-\infty}^{\xi} \frac{\gamma_d}{\rho_d} e^{\lambda_d^-(\xi-t)} \frac{1}{\gamma_d} \left[\beta_d N_d V\left(t\right) + \alpha_d I_d\left(t\right)\right] dt \\
\qquad + \int_{\xi}^{\infty} \frac{\gamma_d}{\rho_d} e^{\lambda_d^+(\xi-t)} \frac{1}{\gamma_d} \left[\beta_d N_d V\left(t\right) + \alpha_d I_d\left(t\right)\right] dt, \\
V\left(\xi\right) = \int_{-\infty}^{\xi} \frac{d_v+d_n}{\rho_v} e^{\lambda_v^-(\xi-t)} \frac{1}{d_v+d_n} H_v(t) dt + \int_{\xi}^{\infty} \frac{d_v+d_n}{\rho_v} e^{\lambda_v^+(\xi-t)} \frac{1}{d_v+d_n} H_v(t) dt.
\end{cases}
$$

Namely,

$$
\begin{cases}
I_c\left(\xi\right) \leq \left(\mathcal{V}^{-1}\mathcal{F}\right)_1 \int_{-\infty}^{\xi} \frac{d_{ic}}{s+v_c} e^{-\frac{d_{ic}}{s+v_c}(\xi-t)} N(t) dt, \\
E_w\left(\xi\right) \leq \left(\mathcal{V}^{-1}\mathcal{F}\right)_2 \left[\int_{-\infty}^{\xi} \frac{\mu_w}{\rho_w} e^{\lambda_w^-(\xi-t)} N(t) dt + \int_{\xi}^{\infty} \frac{\mu_w}{\rho_w} e^{\lambda_w^+(\xi-t)} N(t) dt\right], \\
I_w\left(\xi\right) \leq \left(\mathcal{V}^{-1}\mathcal{F}\right)_3 \left[\int_{-\infty}^{\xi} \frac{d_{iw}}{s} e^{-\frac{d_{iw}}{s}(\xi-t)} \int_{-\infty}^{t} \frac{\mu_w}{\rho_w} e^{\lambda_w^-(t-s)} N(s) ds dt \right. \\
\qquad \left. + \int_{-\infty}^{\xi} \frac{d_{iw}}{s} e^{-\frac{d_{iw}}{s}(\xi-t)} \int_{t}^{\infty} \frac{\mu_w}{\rho_w} e^{\lambda_w^+(t-s)} N(s) ds dt\right], \\
I_d\left(\xi\right) \leq \left(\mathcal{V}^{-1}\mathcal{F}\right)_4 \left[\int_{-\infty}^{\xi} \frac{\gamma_d}{\rho_d} e^{\lambda_d^-(\xi-t)} N(t) dt + \int_{\xi}^{\infty} \frac{\gamma_d}{\rho_d} e^{\lambda_d^+(\xi-t)} N(t) dt\right], \\
V\left(\xi\right) \leq \left(\mathcal{V}^{-1}\mathcal{F}\right)_5 \left[\int_{-\infty}^{\xi} \frac{d_v+d_n}{\rho_v} e^{\lambda_v^-(\xi-t)} N(t) dt + \int_{\xi}^{\infty} \frac{d_v+d_n}{\rho_v} e^{\lambda_v^+(\xi-t)} N(t) dt\right],
\end{cases}
$$

where $\left(\mathcal{V}^{-1}\mathcal{F}\right)_i$ denotes the $i-$th row of the matrix $\mathcal{V}^{-1}\mathcal{F}$ and

$$
\mathcal{V}^{-1}\mathcal{F} = \begin{pmatrix}
\frac{\alpha_c}{d_{ic}} & 0 & 0 & 0 & \frac{\beta_c S_{c0}}{d_{ic}} \\
0 & \frac{\alpha_{ew}}{\mu_w} & \frac{\alpha_{iw}}{\mu_w} & 0 & \frac{\beta_w S_{w0}}{\mu_w} \\
0 & \frac{\alpha_{ew}}{d_{iw}} & \frac{\alpha_{iw}}{d_{iw}} & 0 & \frac{\beta_w S_{w0}}{d_{iw}} \\
0 & 0 & 0 & \frac{\alpha_d}{\gamma_d} & \frac{\beta_d N_d}{\gamma_d} \\
\frac{r_c}{d_v+d_n} & \frac{r_{ew}}{d_v+d_n} & \frac{r_{iw}}{d_v+d_n} & \frac{r_d}{d_v+d_n} & 0
\end{pmatrix}, \quad N(t) = \begin{pmatrix} I_c(t) \\ E_w(t) \\ I_w(t) \\ I_d(t) \\ V(t) \end{pmatrix}.
$$

Let $\sup_{\xi \in \mathbb{R}} I_c\left(\xi\right) = I_c^0$, $\sup_{\xi \in \mathbb{R}} E_w\left(\xi\right) = E_w^0$, $\sup_{\xi \in \mathbb{R}} I_w\left(\xi\right) = I_w^0$, $\sup_{\xi \in \mathbb{R}} I_w\left(\xi\right) = I_d^0$ and $\sup_{\xi \in \mathbb{R}} V\left(\xi\right) = V^0$. Then $N^0 := (I_c^0, E_w^0, I_w^0, I_d^0, V^0)^T \geq \mathbf{0}$ and $N^0 \neq \mathbf{0}$, where T is the transpose. Furthermore, we have

$$
(2.4) \qquad\qquad N^0 \leq \left(\mathcal{V}^{-1}\mathcal{F}\right) N^0.
$$

If $\rho\left(\mathcal{F}\mathcal{V}^{-1}\right) := \rho_0 < 1$, then there exists a nontrivial vector $P := (p_1, p_2, 0, p_4, p_5) \geq 0$ (page 16, Theorem 3.5 of [3]) such that $\left(\mathcal{F}\mathcal{V}^{-1}\right) P = \rho_0 P$. It is easy to verify that $p_1 > 0, p_2 > 0, p_4 > 0, p_5 > 0$. Then $\mathcal{V}^{-1} P > 0$ and there holds $\left(\mathcal{V}^{-1}\mathcal{F}\right)\left(\mathcal{V}^{-1}P\right) = \mathcal{V}^{-1}\left(\mathcal{F}\mathcal{V}^{-1}\right) P = \rho_0 \mathcal{V}^{-1} P$, which implies that $\rho_0 < 1$ is a nonnegative engeinvalue of the matrix $\mathcal{V}^{-1}\mathcal{F}$ with positive engeinvector $\mathcal{V}^{-1}P$. It is easy to show that $\mathcal{V}^{-1}\mathcal{F}$ is irreducible, that is, $\left(\mathcal{V}^{-1}\mathcal{F} + \mathcal{I}\right)^4 > 0$. Then the Perron-Frobenius theorem yields that $\rho\left(\mathcal{V}^{-1}\mathcal{F}\right) = \rho_0 < 1$. So, iterating (2.4) yields $N^0 = \mathbf{0}$, a contradiction. $\qquad\square$

3. Existence of traveling wave solutions

In the following, we prove the existence of traveling waves of (1.1) when $\mathcal{R}_0 > 1$. Linearizing (2.1) for I_c, E_w, I_w, I_d and V in the region $\xi \to -\infty$ where $S_c \to S_{c0}$, $S_w \to S_{w0}$, and setting the remaining variables approaching zero, we have

$$
\begin{cases}
sI_c'(\xi) = \beta_c S_{c0} V + \alpha_c I_c - d_{ic} I_c - v_c I_c'(\xi), \\
sE_w'(\xi) = \beta_w S_{w0} V + \alpha_{ew} E_w + \alpha_{iw} I_w - \mu_w E_w - v_w E_w'(\xi) + D_w E_w''(\xi), \\
sI_w'(\xi) = \mu_w E_w - d_{iw} I_w, \\
sI_d'(\xi) = \beta_d N_d V + \alpha_d I_d - \gamma_d I_d - v_d I_d'(\xi) + D_d I_d''(\xi), \\
sV'(\xi) = r_c I_c + r_{ew} E_w + r_{iw} I_w + r_d I_d - (d_v + d_n) V - v_v V'(\xi) + D_v V''(\xi).
\end{cases}
$$

Looking for the solutions of the form $(I_c, E_w, I_w, I_d, V) = (q_1, q_2, q_3, q_4, q_5) e^{\lambda \xi}$, where $q_i > 0$ and $\lambda > 0$, we have

$$
(3.1) \quad
\begin{cases}
s\lambda q_1 = \beta_c S_{c0} q_5 + \alpha_c q_1 - d_{ic} q_1 - v_c \lambda q_1, \\
s\lambda q_2 = \beta_w S_{w0} q_5 + \alpha_{ew} q_2 + \alpha_{iw} q_3 - \mu_w q_2 - v_w \lambda q_2 + D_w \lambda^2 q_2, \\
s\lambda q_3 = \mu_w q_2 - d_{iw} q_3, \\
s\lambda q_4 = \beta_d N_d q_5 + \alpha_d q_4 - \gamma_d q_4 - v_d \lambda q_4 + D_d \lambda^2 q_4, \\
s\lambda q_5 = r_c q_1 + r_{ew} q_2 + r_{iw} q_3 + r_d q_4 - (d_v + d_n) q_5 - v_v \lambda q_5 + D_v \lambda^2 q_5.
\end{cases}
$$

Let

$$
\widetilde{A} = \begin{pmatrix}
0 & 0 & 0 & 0 & 0 \\
0 & D_w & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & D_d & 0 \\
0 & 0 & 0 & 0 & D_v
\end{pmatrix}, \quad
\widetilde{B} = \begin{pmatrix}
s + v_c & 0 & 0 & 0 & 0 \\
0 & s + v_w & 0 & 0 & 0 \\
0 & 0 & s & 0 & 0 \\
0 & 0 & 0 & s + v_d & 0 \\
0 & 0 & 0 & 0 & s + v_v
\end{pmatrix}
$$

and $M(\lambda, c) := \widetilde{A}\lambda^2 - \widetilde{B}\lambda + \mathcal{F} - \mathcal{V}$. Then (3.1) can be rewritten as $MQ^T = 0$, where $Q = (q_1, q_2, q_3, q_4, q_5)$. Let $A = \mathcal{V}^{-1}\widetilde{A}$ and $B = \mathcal{V}^{-1}\widetilde{B}$. Consider the equation $\left(-A\lambda^2 + B\lambda + I\right)^{-1} \left(\mathcal{V}^{-1}\mathcal{F}\right) Q = Q$. A direct calculation yields

$$
\begin{aligned}
\overline{M}(\lambda, c) : &= \left(-A\lambda^2 + B\lambda + I\right)^{-1} \left(\mathcal{V}^{-1}\mathcal{F}\right) \\
&= \begin{pmatrix}
\frac{\alpha_c}{\Theta_1(\lambda,s)} & 0 & 0 & 0 & \frac{\beta_c S_{c0}}{\Theta_1(\lambda,s)} \\
0 & \frac{\alpha_{ew}}{\Theta_2(\lambda,s)} & \frac{\alpha_{iw}}{\Theta_2(\lambda,s)} & 0 & \frac{\beta_w S_{w0}}{\Theta_2(\lambda,s)} \\
0 & \frac{\alpha_{ew}\mu_w}{\Theta_3(\lambda,s)} & \frac{\alpha_{iw}\mu_w}{\Theta_3(\lambda,s)} & 0 & \frac{\beta_w S_{w0}\mu_w}{\Theta_3(\lambda,s)} \\
0 & 0 & 0 & \frac{\alpha_d}{\Theta_4(\lambda,s)} & \frac{\beta_d N_d}{\Theta_4(\lambda,s)} \\
\frac{r_c}{\Theta_5(\lambda,s)} & \frac{r_{ew}}{\Theta_5(\lambda,s)} & \frac{r_{iw}}{\Theta_5(\lambda,s)} & \frac{r_d}{\Theta_5(\lambda,s)} & 0
\end{pmatrix},
\end{aligned}
$$

where $\Theta_1(\lambda, s) = \lambda(s + v_c) + d_{ic}$, $\Theta_2(\lambda, s) = \lambda(s + v_w) - D_w\lambda^2 + \mu_w$, $\Theta_3(\lambda, s) = (\lambda s + d_{iw})\left[\lambda(s + v_w) - D_w\lambda^2 + \mu_w\right]$, $\Theta_4(\lambda, s) = \lambda(s + v_d) - D_d\lambda^2 + \gamma_d$, $\Theta_5(\lambda, s) = \lambda(s + v_v) - D_v\lambda^2 + (d_v + d_n)$. Take

$$
D = \max\{D_w, D_d, D_v\} \quad \text{and} \quad v_{\min} = \min\{0, v_c, v_w, v_d, v_v\}.
$$

Then for $s \geq -v_{\min}$, we have $s + v_j \geq s + v_{\min}$ with $j = c, w, d, v$.

As $\Theta_i \left(\frac{s+v_{\min}}{2D}, s \right)$ is increasing and nonnegative for $s \in [-v_{\min}, \infty)$, we conclude that the matrix $\overline{M} \left(\frac{s+v_{\min}}{2D}, s \right)$ is decreasing in $s \in [-v_{\min}, +\infty)$. In particular, this matrix becomes $\mathcal{V}^{-1} \mathcal{F}$ when $s = -v_{\min}$ and approaches zero when $s \to +\infty$. Since $\rho(M)$ is continuous and monotonically increasing with respect to the nonnegative matrix M, there exists a unique $s^* > -v_{\min}$ such that $\rho \left(\overline{M} \left(\frac{s^*+v_{\min}}{2D}, s^* \right) \right) = 1$ and $\rho \left(\overline{M} \left(\frac{s+v_{\min}}{2D}, s \right) \right) < 1$ for $s > s^*$. Now we fix $s > s^*$. Since $\Theta_i(\lambda, s)$ is increasing in $\lambda \in \left[0, \frac{s+v_{\min}}{2D} \right]$, we conclude that the matrix $\overline{M}(\lambda, s)$ is decreasing and nonnegative in $\lambda \in \left[0, \frac{s+v_{\min}}{2D} \right]$. Consequently, there exists a $\lambda_s \in \left(0, \frac{s+v_{\min}}{2D} \right)$ such that

$$\rho \left(\overline{M}(\lambda, s) \right) \begin{cases} = 1 & \text{if } \lambda = \lambda_s, \\ < 1 & \text{if } \lambda \in \left(\lambda_s, \frac{s+v_{\min}}{2D} \right], \\ > 1 & \text{if } \lambda \in [0, \lambda_s). \end{cases}$$

Lemma 3.1. *Assume that $\mathcal{R}_0 = \rho \left(\mathcal{F} \mathcal{V}^{-1} \right) > 1$. Then there exists $s^* > -v_{\min}$ such that for each $s > s^*$, there exist $\lambda_s \in \left(0, \frac{s+v_{\min}}{2D} \right)$ and $Q_s > 0$ satisfying $\det M(\lambda_s, s) = 0$ and $M(\lambda_s, s) Q_s = 0$.*

Proof. Following the above arguments, we know that $\rho \left(\overline{M}(\lambda_s, s) \right) = 1$. Then the Perron-Frobenius theorem implies that there exists a $Q_s \in \mathbb{R}^5$ with positive components such that $\overline{M}(\lambda_s, s) Q_s = Q_s$. Multiplying the matrix $-A\lambda_s^2 + B\lambda_s + I$ on the two sides of the last equality, we have $\left(A\lambda_s^2 - B\lambda_s + \mathcal{V}^{-1}\mathcal{F} - I \right) Q_s = 0$. Multiplying $\mathcal{V}$ to the above equality yields $M(\lambda, s)Q_s = 0$, completing the proof. $\square$

In the sequel, we let $Q_s := (q_1, q_2, q_3, q_4, q_5)^T$ as obtained in Lemma 3.1.

Lemma 3.2. *The vector valued map $\Phi(x) = (\phi_1(x), \phi_2(x), \phi_3(x), \phi_4(x), \phi_5(x))^T$ with $\phi_i(x) = q_i e^{\lambda_s x}$ satisfies the following system:*

$$(3.2) \quad s\phi_1'(x) = \beta_c S_{c0} \phi_5(x) + \alpha_c \phi_1(x) - d_{ic}\phi_1(x) - v_c \phi_1'(x),$$

$$(3.3) \quad s\phi_2'(x) = \beta_w S_{w0} \phi_5(x) + \alpha_{ew}\phi_2(x) + \alpha_{iw}\phi_3(x)$$
$$- \mu_w \phi_2(x) - v_w \phi_2'(x) + D_w \phi_2''(x),$$

$$(3.4) \quad s\phi_3'(x) = \mu_w \phi_2(x) - d_{iw}\phi_3(x),$$

$$(3.5) \quad s\phi_4'(x) = \beta_d N_d \phi_5(x) + \alpha_d \phi_4(x) - \gamma_d \phi_4(x) - v_d \phi_4'(x) + D_d \phi_4''(x),$$

$$(3.6) \quad s\phi_5'(x) = r_c \phi_1(x) + r_{ew}\phi_2(x) + r_{iw}\phi_3(x) + r_d \phi_4(x)$$
$$- (d_v + d_n) \phi_5(x) - v_v \phi_5'(x) + D_v \phi_5''(x).$$

Lemma 3.3. *For each $\varpi > 0$ sufficiently small and $\rho > 1$ large enough, the vector valued map $P(x) = (p_1(x), p_2(x))^T$ defined by*

$$p_1(x) = \max\{1 - \rho e^{\varpi x}, 0\} S_{c0} \quad and \quad p_2(x) = \max\{1 - \rho e^{\varpi x}, 0\} S_{w0}$$

satisfies the following system of differential inequalities:

$$(3.7) \quad D_c p_1''(x) - (s + v_c)p_1'(x) - \beta_c p_1(x) \phi_5(x) - \alpha_c \frac{p_1(x) \phi_1(x)}{p_1(x) + \phi_1(x)} \geq 0,$$

$$(3.8) \qquad\qquad D_w p_2''(x) - (s + v_w)p_2'(x) - \beta_w p_2(x) \phi_5(x)$$
$$- \alpha_{ew} \frac{p_2(x) \phi_2(x)}{p_2(x) + \phi_2(x)} - \alpha_{iw} \frac{p_2(x) \phi_3(x)}{p_2(x) + \phi_3(x)} \geq 0$$

for any $x < X' := -\frac{1}{\varpi} \ln \rho$.

Proof. When $x < X'$, $1 - \rho e^{\varpi x} > 0$ and $p_1(x) = S_{c0}(1 - \rho e^{\varpi x})$. One has

$$
\begin{aligned}
& D_c p_1''(x) - (s + v_c)p_1'(x) - \beta_c p_1(x)\phi_5(x) - \alpha_c \frac{p_1(x)\phi_1(x)}{p_1(x) + \phi_1(x)} \\
\geq{} & -\rho\varpi^2 D_c S_{c0} e^{\varpi x} + (s + v_c)\rho\varpi S_{c0} e^{\varpi x} - q_5 \beta_c S_{c0}(1 - \rho e^{\varpi x})e^{\lambda_s x} - \alpha_c q_1 e^{\lambda_s x} \\
\geq{} & \left[\rho\varpi S_{c0}((s + v_c) - \varpi D_c) - (\alpha_c q_1 + q_5 \beta_c S_{c0})e^{-(\lambda_s - \varpi)\frac{1}{\varpi}\ln\rho}\right]e^{\varpi x} \\
={} & \left[\rho\varpi S_{c0}((s + v_c) - \varpi D_c) - (\alpha_c q_1 + q_5 \beta_c S_{c0})\rho^{-(\lambda_s - \varpi)\frac{1}{\varpi}}\right]e^{\varpi x}.
\end{aligned}
$$

Keeping $\rho\varpi = 1$ and letting $\rho \to +\infty$, there exists $\rho > 0$ and $\varpi > 0$ such that $\rho\varpi S_{c0}((s + v_c) - \varpi D_c) - (\alpha_c q_1 + q_5 \beta_c S_{c0})\rho^{-(\lambda_s - \varpi)\frac{1}{\varpi}} > 0$, which implies that (3.7) holds. Similarly, we can prove (3.8). This completes the proof. $\quad\square$

Lemma 3.4. *Let $\varepsilon > 0$ be small enough with $\varepsilon < \frac{\varpi}{2}$, $\varepsilon < \frac{\lambda_s}{2}$ and $\lambda_s + \varepsilon < \frac{s + v_{\min}}{2D}$. Then the function*

$$
\Psi(x) = (\psi_1(x), \psi_2(x), \psi_3(x), \psi_4(x), \psi_5(x))^T = Q_s e^{\lambda_s x}\max\{1 - Me^{\varepsilon x}, 0\}
$$

satisfies the following inequalities:

$$
\text{(3.9)} \qquad (s + v_c)\psi_1'(x) \leq \beta_c p_1(x)\psi_5(x) + \alpha_c \frac{p_1(x)\psi_1(x)}{p_1(x) + \psi_1(x)} - d_{ic}\psi_1(x),
$$

$$
\text{(3.10)} \qquad (s + v_w)\psi_2'(x) \leq \beta_w p_2(x)\psi_5(x) + \alpha_{ew}\frac{p_2(x)\psi_2(x)}{p_2(x) + \psi_2(x) + \phi_3(x)}
$$
$$
+ \alpha_{iw}\frac{p_2(x)\psi_3(x)}{p_2(x) + \phi_2(x) + \psi_3(x)} - \mu_w\psi_2(x) + D_w\psi_2'',
$$

$$
\text{(3.11)} \qquad s\psi_3'(x) \leq \mu_w\psi_2(x) - d_{iw}\psi_3(x),
$$

$$
\text{(3.12)} \qquad (s + v_d)\psi_4'(x) \leq \beta_d N_d\psi_5(x) - \beta_d\psi_4(x)\phi_5(x)
$$
$$
+ \alpha_d\frac{(N_d - \psi_4(x))\psi_4(x)}{N_d} - \gamma_d\psi_4(x) + D_d\psi_4'',
$$

$$
\text{(3.13)} \qquad (s + v_v)\psi_5'(x) \leq r_c\psi_1(x) + r_{ew}\psi_2(x) + r_{iw}\psi_3(x) + r_d\psi_4(x)
$$
$$
- (d_v + d_n)\psi_5(x) + D_v\psi_5''
$$

for $x < X'' := -\frac{1}{\varepsilon}\ln M$, where $M > 0$ is sufficiently large so that $X'' < X'$.

Proof. When $x < X'' < X'$, $\psi_i(x) = q_i e^{\lambda_s x}(1 - Me^{\varepsilon x})$, $p_1(x) = S_{c0}(1 - \rho e^{\varpi x})$ and $p_2(x) = S_{w0}(1 - \rho e^{\varpi x})$, where $i = 1, 2, 3, 4, 5$. Consequently, we have

$$s\psi_1'(x) - \beta_c p_1(x)\psi_5(x) - \alpha_c \frac{p_1(x)\psi_1(x)}{p_1(x) + \psi_1(x)} + d_{ic}\psi_1(x) + v_c\psi_1'(x)$$

$$= -(s + v_c)q_1(\lambda_s + \varepsilon)Me^{(\lambda_s + \varepsilon)x} + q_5\beta_c S_{c0}e^{\lambda_s x}\left(Me^{\varepsilon x} + \rho e^{\varpi x} - \rho Me^{(\varpi + \varepsilon)x}\right)$$

$$+ \alpha_c q_1 Me^{(\lambda_s + \varepsilon)x} + \frac{q_1^2 \alpha_c e^{2\lambda_s x}(1 - Me^{\varepsilon x})^2}{S_{c0}(1 - \rho e^{\varpi x}) + q_1 e^{\lambda_s x}(1 - Me^{\varepsilon x})} - q_1 d_{ic}Me^{(\lambda_s + \varepsilon)x}$$

$$\leq M\left[-(s + v_c)(\lambda_s + \varepsilon)q_1 + \alpha_c q_1 - d_{ic}q_1 + q_5\beta_c S_{c0}\right]e^{(\lambda_s + \varepsilon)x}$$

$$+ \frac{q_1^2 \alpha_c e^{2\lambda_s x}}{S_{c0}(1 - \rho e^{\varpi x})} + \rho q_5 \beta_c S_{c0}e^{(\lambda_s + \varpi)x}$$

$$\leq \left\{-\varepsilon q_1 M + \frac{q_1^2 \alpha_c}{S_{c0}\left(1 - \rho e^{-\varpi \frac{1}{\varepsilon}\ln M}\right)} + \rho q_5 \beta_c S_{c0}e^{-(\varpi - \varepsilon)\frac{1}{\varepsilon}\ln M}\right\}e^{(\lambda_s + \varepsilon)x}.$$

Then for sufficiently large $M > 0$, we have that (3.9) holds.

The proofs for (3.10)-(3.13) are similar and thus are omitted. $\qquad\square$

Let $X^* := -\frac{1}{\varepsilon}\ln\frac{M(\lambda_s + \varepsilon)}{\lambda_s} < X''$. It is obvious that $\varphi_i(\cdot)$ is increasing on $(-\infty, X^*]$. For $X > -X^*$, we define

$$\Gamma = \left\{ \begin{pmatrix} \chi_1(\cdot) \\ \chi_2(\cdot) \\ \varphi_1(\cdot) \\ \varphi_2(\cdot) \\ \varphi_3(\cdot) \\ \varphi_4(\cdot) \\ \varphi_5(\cdot) \end{pmatrix} \in C\left(\bar{\Omega}, \mathbb{R}^7\right) \middle| \begin{array}{l} p_1(x) \leq \chi_1(x) \leq S_{c0}, p_2(x) \leq \chi_2(x) \leq S_{w0} \\ \text{for } x \in [-X, X], \chi_i(\pm X) = p_i(\pm X), \\ \psi_j(x) \leq \varphi_j(x) \leq \phi_j(x) \text{ for } x \in [-X, X], \\ i = 1, 2 \text{ and } j = 1, 2, 3, 4, 5; \\ \varphi_1(-X) = \psi_1(-X), \varphi_2(\pm X) = \psi_2(\pm X), \\ \varphi_3(-X) = \psi_3(-X), \varphi_4(\pm X) = \psi_4(\pm X), \\ \varphi_5(\pm X) = \psi_5(\pm X). \\ \text{Furthermore, } \varphi_4(x) \leq N_d \text{ for } x \in [-X, X]. \end{array} \right\},$$

where $\bar{\Omega} = [-X, X]$. For any given $(\chi_1(\cdot), \chi_2(\cdot), \varphi_1(\cdot), \varphi_2(\cdot), \varphi_3(\cdot), \varphi_4(\cdot), \varphi_5(\cdot)) \in \Gamma$, we consider the following boundary value problems:

$$(3.14) \qquad -D_c S_c''(x) + (s + v_c)S_c'(x) + (\beta_c\varphi_5(x) + \alpha_c)S_c(x) = \alpha_c g_1(x),$$

$$(3.15) \qquad -D_w S_w''(x) + (s + v_w)S_w'(x) + (\beta_w\varphi_5(x) + \alpha_{ew} + \alpha_{iw})S_w(x)$$
$$= \alpha_{ew}g_{21}(x) + \alpha_{iw}g_{22}(x),$$

$$(3.16) \qquad (s + v_c)I_c'(x) + d_{ic}I_c(x) = \beta_c\chi_1(x)\varphi_5(x) + \alpha_c f_1(x),$$

$$(3.17) \qquad -D_w E_w''(x) + (s + v_w)E_w'(x) + \mu_w E_w(x)$$
$$= \beta_c\chi_2(x)\varphi_5(x) + \alpha_{ew}f_{21}(x) + \alpha_{iw}f_{22}(x),$$

$$(3.18) \qquad sI_w'(x) + d_{iw}I_w(x) = \mu_w\varphi_2(x),$$

$$(3.19) \qquad -D_d I_d''(x) + (s + v_d)I_d'(x) + (\gamma_d + \alpha_d + \beta_d\varphi_5(x))I_d(x)$$
$$= \beta_d N_d\varphi_5(x) + \frac{\alpha_d}{N_d}(2N_d - \varphi_4(x))\varphi_4(x),$$

$$(3.20) \qquad -D_v V''(x) + (s + v_v)V'(x) + (d_v + d_n)V(x)$$
$$= r_c\varphi_1(x) + r_{ew}\varphi_2(x) + r_{iw}\varphi_3(x) + r_d\varphi_4(x),$$

with

$$
\begin{aligned}
S_c(\pm X) &= p_1(\pm X), S_w(\pm X) = p_2(\pm X), I_c(-X) = \psi_1(-X), \\
(3.21)\quad E_w(\pm X) &= \psi_2(\pm X), I_w(-X) = \psi_3(-X), \\
I_d(\pm X) &= \psi_4(\pm X), V(\pm X) = \psi_5(\pm X),
\end{aligned}
$$

where

$$
f_1[\chi_1,\varphi_1](x) = \begin{cases} \frac{\alpha_c\chi_1(x)\varphi_1(x)}{\chi_1(x)+\varphi_1(x)}, & \chi_1(x)\varphi_1(x) \neq 0, \\ 0, & \chi_1(x)\varphi_1(x) = 0, \end{cases}
$$

$$
g_1[\chi_1,\varphi_1](x) = \begin{cases} \frac{\alpha_c\chi_1^2(x)}{\chi_1(x)+\varphi_1(x)}, & \chi_1(x) \neq 0, \\ 0, & \chi_1(x) = 0, \end{cases}
$$

$$
f_{21}[\chi_2,\varphi_2,\varphi_3](x) = \begin{cases} \frac{\alpha_{ew}\chi_2(x)\varphi_2(x)}{\chi_2(x)+\varphi_2(x)+\varphi_3(x)}, & \chi_2(x)\varphi_2(x) \neq 0, \\ 0, & \chi_2(x)\varphi_2(x) = 0, \end{cases}
$$

$$
f_{22}[\chi_2,\varphi_2,\varphi_3](x) = \begin{cases} \frac{\alpha_{iw}\chi_2(x)\varphi_3(x)}{\chi_2(x)+\varphi_2(x)+\varphi_3(x)}, & \chi_2(x)\varphi_3(x) \neq 0, \\ 0, & \chi_2(x)\varphi_3(x) = 0, \end{cases}
$$

$$
g_{21}[\chi_2,\varphi_2,\varphi_3](x) = \begin{cases} \frac{\alpha_{ew}\chi_2(x)(\chi_2(x)+\varphi_3(x))}{\chi_2(x)+\varphi_2(x)+\varphi_3(x)}, & \chi_2(x)(\chi_2(x)+\varphi_3(x)) \neq 0, \\ 0, & \chi_2(x)(\chi_2(x)+\varphi_3(x)) = 0, \end{cases}
$$

$$
g_{22}[\chi_2,\varphi_2,\varphi_3](x) = \begin{cases} \frac{\alpha_{iw}\chi_2(x)(\chi_2(x)+\varphi_2(x))}{\chi_2(x)+\varphi_2(x)+\varphi_3(x)}, & \chi_2(x)(\chi_2(x)+\varphi_2(x)) \neq 0, \\ 0, & \chi_2(x)(\chi_2(x)+\varphi_2(x)) = 0. \end{cases}
$$

It is not difficult to verify that $f_1(x), f_{21}(x), f_{22}(x), g_1(x), g_{21}(x), g_{22}(x)$ are continuous functions of $x \in [-X, X]$. Then the problems (3.14)-(3.20) and (3.21) admit a unique solution $(S_c(\cdot), S_w(\cdot), I_c(\cdot), E_w(\cdot), I_w(\cdot), I_d(\cdot), V(\cdot))$ with $I_c, I_w \in C^1[-X, X]$, $E_w, V \in C^2[-X, X]$ and $S_c, S_w, I_d \in W^{2,p}(-X, X)$, where $p \geq 2$. (In fact, (3.14), (3.15) and (3.19) hold for almost everywhere $x \in (-X, X)$; see [1, 2, 6].) This then gives an operator $T = (T_1, T_2, T_3, T_4, T_5, T_6, T_7)$ defined on Γ as

$$
\begin{aligned}
S_c &= T_1(\chi_1, \chi_2, \varphi_1, \varphi_2, \varphi_3, \varphi_4, \varphi_5), \quad S_w = T_2(\chi_1, \chi_2, \varphi_1, \varphi_2, \varphi_3, \varphi_4, \varphi_5), \\
I_c &= T_3(\chi_1, \chi_2, \varphi_1, \varphi_2, \varphi_3, \varphi_4, \varphi_5), \quad E_w = T_4(\chi_1, \chi_2, \varphi_1, \varphi_2, \varphi_3, \varphi_4, \varphi_5), \\
I_w &= T_5(\chi_1, \chi_2, \varphi_1, \varphi_2, \varphi_3, \varphi_4, \varphi_5), \quad I_d = T_6(\chi_1, \chi_2, \varphi_1, \varphi_2, \varphi_3, \varphi_4, \varphi_5), \\
V &= T_7(\chi_1, \chi_2, \varphi_1, \varphi_2, \varphi_3, \varphi_4, \varphi_5).
\end{aligned}
$$

Theorem 3.5. *The operator T maps Γ into Γ.*

Proof. Let $(\chi_1(\cdot), \chi_2(\cdot), \varphi_1(\cdot), \varphi_2(\cdot), \varphi_3(\cdot), \varphi_4(\cdot), \varphi_5(\cdot)) \in \Gamma$ and

$$
(S_c(\cdot), S_w(\cdot), I_c(\cdot), E_w(\cdot), I_w(\cdot), I_d(\cdot), V(\cdot)) = T(\chi_1, \chi_2, \varphi_1, \varphi_2, \varphi_3, \varphi_4, \varphi_5)(\cdot).
$$

By virtue of the embedding theorem, we have $S_c(\cdot), S_w(\cdot), I_d(\cdot) \in C([-X, X], \mathbb{R})$. Now we show that $p_2(x) \leq S_w(x) \leq S_{w0}$ for $x \in [-X, X]$. Since 0 is a subsolution of (3.15), by the maximum principle [6] we have $S_w(x) \geq 0$ for $x \in [-X, X]$. Furthermore, since $(\beta_w\varphi_5(x) + \alpha_{ew} + \alpha_{iw})S_{w0} \geq \alpha_{ew}g_{21}(x) + \alpha_{iw}g_{22}(x)$ for $x \in (-X, X)$ and $S_w(\pm X) \leq S_{w0}$, S_{w0} is a supersolution of (3.15), and hence we have

$S_w(x) \leq S_{w0}$ for $x \in [-X, X]$. Note that $S_w(-X) = p_2(-X)$ and $S_w(X') \geq p_2(X') = 0$. Then for $x \in (-X, X')$, by (3.8) we have

$$
\begin{aligned}
0 &\geq -D_w p_2''(x) + (s + v_w) p_2'(x) + \beta_w \phi_5(x) p_2(x) \\
&\quad + \alpha_{ew} \frac{p_2(x) \phi_2(x)}{p_2(x) + \phi_2(x)} + \alpha_{iw} \frac{p_2(x) \phi_3(x)}{p_2(x) + \phi_3(x)} \\
&\geq -D_w p_2''(x) + (s + v_w) p_2'(x) + \beta_w \varphi_5(x) p_2(x) \\
&\quad + \alpha_{ew} \frac{p_2(x) \varphi_2(x)}{p_2(x) + \varphi_2(x) + \varphi_3(x)} + \alpha_{iw} \frac{p_2(x) \varphi_3(x)}{p_2(x) + \varphi_2(x) + \varphi_3(x)} \\
&= -D_w p_2''(x) + (s + v_w) p_2'(x) + (\beta_w \varphi_5(x) + \alpha_{ew} + \alpha_{iw}) p_2(x) \\
&\quad - \alpha_{ew} \frac{p_2(x)(p_2(x) + \varphi_3(x))}{p_2(x) + \varphi_2(x) + \varphi_3(x)} - \alpha_{iw} \frac{p_2(x)(p_2(x) + \varphi_2(x))}{p_2(x) + \varphi_2(x) + \varphi_3(x)} \\
&\geq -D_w p_2''(x) + (s + v_w) p_2'(x) + (\beta_w \phi_5(x) + \alpha_{ew} + \alpha_{iw}) p_2(x) \\
&\quad - \alpha_{ew} g_{21}(x) - \alpha_{iw} g_{22}(x),
\end{aligned}
$$

which implies that $p_2(\cdot)$ is a subsolution of (3.15) on $[-X, X']$. Here we used the fact that the function $\frac{u(u+a)}{u+v+a}$ is nondecreasing on u and nonincreasing on v for $(u, v) \in (0, +\infty) \times [0, +\infty)$, where $a \geq 0$. Consequently, the maximum principle yields that $S_w(x) \geq p_2(x)$ for $x \in [-X, X']$. Combining the above arguments, we have $p_2(x) \leq S_w(x) \leq S_{w0}$ for $x \in [-X, X]$. Similarly, we can confirm that $p_1(x) \leq S_c(x) \leq S_{c0}$ for $x \in [-X, X]$.

We now show that $\psi_2(x) \leq E_w(x) \leq \phi_2(x)$ for $x \in [-X, X]$. By (3.3), we have

$$
\begin{aligned}
&-D_w \phi_2''(x) + (s + v_w) \phi_2'(x) + \mu_w \phi_2(x) \\
&= \beta_w S_{w0} \phi_5(x) + \alpha_{ew} \phi_2(x) + \alpha_{iw} \phi_3(x) \\
&\geq \beta_c \chi_2(x) \varphi_5(x) + \frac{\alpha_{ew} \chi_2(x) \phi_2(x)}{\phi_2(x) + \chi_2(x)} + \frac{\alpha_{ew} \chi_2(x) \phi_3(x)}{\phi_3(x) + \chi_2(x)} \\
&\geq \beta_c \chi_2(x) \varphi_5(x) + \alpha_{ew} f_{21}(x) + \alpha_{iw} f_{22}(x).
\end{aligned}
$$

This, combined with $\psi_2(\pm X) = E_w(\pm X) \leq \phi_2(\pm X)$, implies that $\phi_2(\cdot)$ is a supersolution of (3.17) on $[-X, X]$. The maximum principle implies that $E_w(x) \leq \phi_2(x)$ for $x \in [-X, X]$. Furthermore, since 0 is a subsolution of (3.17), we have that $E_w(x) \geq 0$ for $x \in [-X, X]$. Then for $x \in (-X, X'')$, by (3.10) we have

$$
\begin{aligned}
&-D_w \psi_2'' + (s + v_w) \psi_2'(x) + \mu_w \psi_2(x) \\
&\leq \beta_w p_2(x) \psi_5(x) + \alpha_{ew} \frac{p_2(x) \psi_2(x)}{p_2(x) + \psi_2(x) + \phi_3(x)} + \alpha_{iw} \frac{p_2(x) \psi_3(x)}{p_2(x) + \phi_2(x) + \psi_3(x)} \\
&\leq \beta_w \chi_2(x) \varphi_5(x) + \alpha_{ew} \frac{\chi_2(x) \varphi_2(x)}{\chi_2(x) + \varphi_2(x) + \varphi_3(x)} + \alpha_{iw} \frac{\chi_2(x) \varphi_3(x)}{\chi_2(x) + \varphi_2(x) + \varphi_3(x)} \\
&= \beta_c \chi_2(x) \varphi_5(x) + \alpha_{ew} f_{21}(x) + \alpha_{iw} f_{22}(x).
\end{aligned}
$$

In view of $\psi_2(-X) = E_w(-X)$ and $0 \leq \psi_2(X'') \leq E_w(X'')$, an application of the maximum principle yields $E_w(x) \geq \psi_2(x)$ for $x \in [-X, X'']$. Thus, we have proved that $\psi_2(x) \leq E_w(x) \leq \phi_2(x)$ for $x \in [-X, X]$. Similarly, we can prove that $\psi_1(x) \leq I_c(x) \leq \phi_1(x)$, $\psi_3(x) \leq I_w(x) \leq \phi_3(x)$ and $\psi_5(x) \leq V(x) \leq \phi_5(x)$ for $x \in [-X, X]$.

Finally, we prove that $\psi_4(x) \leq I_d(x) \leq \min\{\phi_4(x), N_d\}$ for any $x \in [-X, X]$. Note that $2N_d x - x^2$ is increasing on $x \in (0, N_d)$. Then we have

$$(\gamma_d + \alpha_d + \beta_d \varphi_5(x)) N_d > \beta_d N_d \varphi_5(x) + \frac{\alpha_d}{N_d}(2N_d - \varphi_4(x)) \varphi_4(x),$$

which implies that N_d is a supersolution of (3.19) on $x \in [-X, X]$. Consequently, we have $I_d(x) \leq N_d$ for $x \in [-X, X]$. Similarly, we have $I_d(x) \geq 0$ for $x \in [-X, X]$. Furthermore, by (3.12) we have

$$-D_d \psi_4'' + (s + v_d) \psi_4'(x) + (\gamma_d + \alpha_d + \beta_d \varphi_5(x)) \psi_4(x)$$
$$\leq -D_d \psi_4'' + (s + v_d) \psi_4'(x) + (\gamma_d + \alpha_d + \beta_d \phi_5(x)) \psi_4(x)$$
$$\leq \beta_d N_d \psi_5(x) + \alpha_d \frac{(2N_d - \psi_4(x)) \psi_4(x)}{N_d}$$
$$\leq \beta_d N_d \varphi_5(x) + \alpha_d \frac{(2N_d - \varphi_4(x)) \varphi_4(x)}{N_d},$$

which implies that $\psi_4(x)$ is a subsolution of (3.19) on $[-X, X'']$. Therefore, we have $I_d(x) \geq \psi_4(x)$ for $x \in [-X, X'']$. On the other hand, inequality (3.5) implies that ϕ_4 is a supersolution of (3.19), and hence $I_d(x) \leq \phi_4(x)$ for $x \in [-X, X]$. $\square$

By the classical embedding theorems, we have that T is a compact operator from Γ into Γ. Now we show that $T : \Gamma \to \Gamma$ is continuous. First, we show that $f_1, f_{21}, f_{22}, g_1, g_{21}, g_{22}$ are continuous operators from Γ to $C([-X, X], \mathbb{R}_+)$. Consider f_1 first. Let $\chi_1^1, \chi_1^2, \varphi_1^1, \varphi_1^2 \in C[-X, X]$ with $p_1(x) \leq \chi_1^1(x), \chi_1^2(x) \leq S_{c0}$ and $\psi_1(x) \leq \varphi_1^1(x), \varphi_1^2(x) \leq \phi_1(x)$ for $x \in [-X, X]$. When $\chi_1^1(x) \varphi_1^1(x) \chi_1^2(x) \varphi_1^2(x) \neq 0$, we have

$$\left| f_1\left[\chi_1^1, \varphi_1^1\right](x) - f_1\left[\chi_1^2, \varphi_1^2\right](x) \right|$$
$$\leq \frac{\alpha_c \chi_1^1(x) \chi_1^2(x) \left|\varphi_1^1(x) - \varphi_1^2(x)\right|}{\left[\chi_1^1(x) + \varphi_1^1(x)\right]\left[\chi_1^2(x) + \varphi_1^2(x)\right]} + \frac{\alpha_c \varphi_1^1(x) \varphi_1^2(x) \left|\chi_1^1(x) - \chi_1^2(x)\right|}{\left[\chi_1^1(x) + \varphi_1^1(x)\right]\left[\chi_1^2(x) + \varphi_1^2(x)\right]}$$
$$\leq \alpha_c \left|\varphi_1^1(x) - \varphi_1^2(x)\right| + \alpha_c \left|\chi_1^1(x) - \chi_1^2(x)\right|.$$

When $\chi_1^1(x) \varphi_1^1(x) \chi_1^2(x) \varphi_1^2(x) = 0$ and $\chi_1^1(x) \varphi_1^1(x) + \chi_1^2(x) \varphi_1^2(x) \neq 0$, for example, $\chi_1^1(x) = 0$ and $\chi_1^2(x) \varphi_1^2(x) \neq 0$, we have

$$\left| f_1\left[\chi_1^1, \varphi_1^1\right](x) - f_1\left[\chi_1^2, \varphi_1^2\right](x) \right| = \left| \frac{\alpha_c \varphi_1^2(x)}{\chi_1^2(x) + \varphi_1^2(x)} \right| \left|\chi_1^1(x) - \chi_1^2(x)\right|$$
$$\leq \alpha_c \left|\chi_1^1(x) - \chi_1^2(x)\right|.$$

Therefore, f_1 is continuous. Similarly, we can prove that $f_{21}, f_{22}, g_1, g_{21}, g_{22}$ are continuous. Consequently, using the continuous dependence of solutions of ODEs on initial values and the standard elliptic estimates [1, 2, 6], we have that T_3, T_4, T_5, T_7 are continuous operators on Γ. Furthermore, let $S_c^i = T_1\left(\chi_1^i, \chi_2^i, \varphi_1^i, \varphi_2^i, \varphi_3^i, \varphi_4^i, \varphi_5^i\right)$, $i = 1, 2$, where $\left(\chi_1^i, \chi_2^i, \varphi_1^i, \varphi_2^i, \varphi_3^i, \varphi_4^i, \varphi_5^i\right) \in \Gamma$. Then, we have

$$-D_c\left(S_c^1 - S_c^2\right)''(x) + (s + v_c)\left(S_c^1 - S_c^2\right)'(x) + \left(\alpha_c + \beta_c \varphi_5^1(x)\right)\left(S_c^1 - S_c^2\right)(x)$$
$$= \alpha_c\left[g_1\left[\chi_1^1, \varphi_1^1\right](x) - g_1\left[\chi_1^2, \varphi_1^2\right](x)\right] + \beta_c S_c^2(x)\left[\varphi_5^1(x) - \varphi_5^2(x)\right].$$

Note that $0 \leq S_c^2(x) \leq S_{c0}$ for $x \in [-X, X]$. Applying the standard elliptic estimates to the last equality yields that T_1 is a continuous operator on Γ. Similarly, T_2 and T_6 are also continuous operators on Γ.

Combining the above arguments, we know that $T : \Gamma \to \Gamma$ is a completely continuous operator. Hence, Schauder's fixed point theorem implies that there exists $(S_{c,X}, S_{w,X}, I_{c,X}, E_{w,X}, I_{w,X}, I_{d,X} V_X) \in \Gamma$ such that

$$
\begin{aligned}
& (S_{c,X}(x), S_{w,X}(x), I_{c,X}(x), E_{w,X}(x), I_{w,X}(x), I_{d,X}(x), V_X(x)) \\
& \quad = T(S_{c,X}, S_{w,X}, I_{c,X}, E_{w,X}, I_{w,X}, I_{d,X}, V_X)(x)
\end{aligned}
$$

for $x \in [-X, X]$.

Theorem 3.6. (i) $I_{c,X}, I_{w,X} \in C^2[-X, X]$, $S_{c,X}, S_{w,X}, E_{w,X}, I_{d,X}, V_X \in C^3[-X, X]$.
 (ii) *There exists $M_0 > 0$ such that for any $X > -X^*$, $\|I_{c,X}\|_{C^1(\bar{\Omega})}$, $\|I_{w,X}\|_{C^2(\bar{\Omega})}$,*
$\|V_X\|_{C(\bar{\Omega})}$, $\|S_{c,X}\|_{C^2(\bar{\Omega})}$, $\|S_{w,X}\|_{C^2(\bar{\Omega})}$, $\|E_{w,X}\|_{C^2(\bar{\Omega})} \le M_0$, where $\bar{\Omega} = [-X, X]$.
 (iii) *For each $Y > -X^*$, there exists $C_0 = C_0(Y) > 0$ such that*
(3.22)
$$
\|I_{c,X}''\|_{C(\bar{U})}, \|S_{c,X}'''\|_{C(\bar{U})}, \|S_{w,X}'''\|_{C(\bar{U})}, \|E_{w,X}'''\|_{C(\bar{U})}, \|I_{d,X}\|_{C^3(\bar{U})}, \|V_X\|_{C^3(\bar{U})} \le C_0
$$

for any $X > Y$, where $\bar{U} = [-Y, Y]$.

Proof. (i) First, we have $I_{c,X}, I_{w,X} \in C^1[-X, X]$, $E_{w,X}, V_X \in C^2[-X, X]$ and $S_{c,X}, S_{w,X}, I_{d,X} \in W^{2,p}(-X, X)$, $p \ge 2$. Then by the embedding theorem, we have $S_{c,X}, S_{w,X}, I_{d,X} \in W^{2,p}(-X, X) \hookrightarrow C^{1+\alpha}[-X, X]$ for some $\alpha \in (0, 1)$. Furthermore, it is not difficult to prove that $g_1[S_{c,X}, I_{c,X}](\cdot)$, $g_{21}[S_{w,X}, E_{w,X}, I_{w,X}](\cdot)$, $g_{22}[S_{w,X}, E_{w,X}, I_{w,X}](\cdot) \in C^\alpha[-X, X]$, which further yields $S_{c,X}, S_{w,X}, I_{d,X} \in C^{2,\alpha}[-X, X]$. Then, we can verify that the operators $f_1, f_{21}, f_{22}, g_1, g_{21}$ and g_{22} map the solutions $(S_{c,X}, S_{w,X}, I_{c,X}, E_{w,X}, I_{w,X}, I_{d,X}, V_X)$ into $C^1[-X, X]$. Therefore, $I_{c,X}, I_{w,X} \in C^2[-X, X]$, $S_{c,X}, S_{w,X}, E_{w,X}, I_{d,X}, V_X \in C^3[-X, X]$.
 (ii) It is obvious that $(S_{c,X}, S_{w,X}, I_{c,X}, E_{w,X}, I_{w,X}, I_{d,X}, V_X) \in \Gamma$ satisfies

$$
\begin{aligned}
(3.23) \quad (s + v_c) S_{c,X}' &= -\beta_c S_{c,X} V_X - \alpha_c S_{c,X} I_{c,X} / N_{c,X} + D_c S_{c,X}'', \\
(3.24) \quad (s + v_w) S_{w,X}' &= -\beta_w S_{w,X} V_X - \alpha_{ew} S_{w,X} E_{w,X} / N_{w,X} \\
& \quad - \alpha_{iw} S_{w,X} I_{w,X} / N_{w,X} + D_w S_{w,X}'', \\
(3.25) \quad (s + v_c) I_{c,X}' &= \beta_c S_{c,X} V_X + \alpha_c S_{c,X} I_{c,X} / N_{c,X} - d_{ic} I_{c,X}, \\
(3.26) \quad (s + v_w) E_{w,X}' &= \beta_w S_{w,X} V_X + S_{w,X}[\alpha_{ew} E_{w,X} + \alpha_{iw} I_{w,X}] / N_{w,X} \\
& \quad - \mu_w E_{w,X} + D_w E_{w,X}'', \\
(3.27) \quad s I_{w,X}' &= \mu_w E_{w,X} - d_{iw} I_{w,X}, \\
(3.28) \quad (s + v_d) I_{d,X}' &= \beta_d(N_d - I_{d,X}) V_X + \alpha_d(N_d - I_{d,X}) I_{d,X} / N_d \\
& \quad - \gamma_d I_{d,X} + D_d I_{d,X}'', \\
(3.29) \quad (s + v_v) V_X' &= r_c I_{c,X} + r_{ew} E_{w,X} + r_{iw} I_{w,X} + r_d I_{d,X} \\
& \quad - (d_v + d_n) V_X + D_v V_X'',
\end{aligned}
$$

where $N_{c,X} = S_{c,X} + I_{c,X}$ and $N_{w,X} = S_{w,X} + E_{w,X} + I_{w,X}$.
 Following (3.23), we have

$$
\left[e^{-\frac{s + v_c}{D_c} x} S_{c,X}'(x) \right]' = \frac{1}{D_c} e^{-\frac{s + v_c}{D_c} x} \left[\beta_c S_{c,X}(x) V_X(x) + \alpha_c \frac{S_{c,X}(x) I_{c,X}(x)}{N_{c,X}(x)} \right].
$$

Therefore, for $x \in [-X, X)$ we have

$$S'_{c,X}(x) - e^{\frac{s+v_c}{D_c}(x-X)} S'_{c,X}(X)$$
$$= -\frac{1}{D_c} \int_x^X e^{\frac{s+v_c}{D_c}(x-z)} [\beta_c S_{c,X}(z) V_X(z) + \alpha_c S_{c,X}(z) I_{c,X}(z)/N_{c,X}(z)] dz.$$

Since $S'_{c,X}(X) \leq 0$, we have $S'_{c,X}(x) \leq 0$ for any $x \in [-X, X]$. From (3.23) and (3.25), we further have $(s+v_c) [S'_{c,X}(x) + I'_{c,X}(x)] + d_{ic} I_{c,X}(x) = D_c S''_{c,X}(x)$. Integrating two sides of the above equality from $-X$ to X, we have

$$D_c [S'_{c,X}(X) - S'_{c,X}(-X)] - (s+v_c) [S_{c,X}(X) - S_{c,X}(-X)]$$
$$= (s+v_c) [I_{c,X}(X) - I_{c,X}(-X)] + d_{ic} \int_{-X}^X I_{c,X}(x) dx.$$

Since $S_{c,X}(X) = 0, I_{c,X}(-X) = \psi_1(-X)$ and $S'_{c,X}(-X) \geq p'_1(-X)$, we have

$$\int_{-X}^X I_{c,X}(x) dx \leq \frac{1}{d_{ic}} [-D_c p'_1(-X) + (s+v_c) \psi_1(-X) + (s+v_c) S_{c0}] \leq M_0$$

and $I_{c,X}(X) \leq M_0$ for some $M_0 > 0$, which is independent of $X > -X^*$. Consequently, from (3.25) we get

$$\int_{-X}^X [\beta_c S_{c,X}(z) V_X(z) + \alpha_c S_{c,X}(z) I_{c,X}(z)/N_{c,X}(z)] dz$$
$$\leq (s+v_c) [I_{c,X}(X) - \psi_1(-X)] + d_{ic} \int_{-X}^X I_{c,X}(x) dx \leq M_0$$

for some $M_0 > 0$, which is independent of $X > -X^*$. Integrating two sides of (3.23) from $-X$ to x, we then have

$$D_c S'_{c,X}(x) = D_c S'_{c,X}(-X) + (s+v_c) [S_{c,X}(x) - S_{c,X}(-X)]$$
$$+ \int_{-X}^X [\beta_c S_{c,X}(z) V_X(z) + \alpha_c S_{c,X}(z) I_{c,X}(z)/N_{c,X}(z)] dz$$
$$\geq D_c p'_1(-X) - (s+v_c) S_{c0}.$$

Therefore, there exists $M_0 > 0$ independent of $X > -X^*$ such that $|S'_{c,X}(x)| \leq M_0$ for $x \in [-X, X]$. Similarly, integrating two sides of (3.25) from $-X$ to x, we can find an $M_0 > 0$ independent of $X > -X^*$ such that $|I_{c,X}(x)| \leq M_0$ for $x \in [-X, X]$.

Consider $S_{w,X}, E_{w,X}$ and $I_{w,X}$. From (3.24) and (3.26) we have
(3.30)
$$D_w [S_{w,X}(x) + E_{w,X}(x)]'' - (s+v_w) [S_{w,X}(x) + E_{w,X}(x)]' = \mu_w E_{w,X}(x).$$

Since $S'_{w,X}(X) \leq 0$ and $E'_{w,X}(X) \leq 0$, we can obtain that $S'_{w,X}(x) + E'_{w,X}(x) \leq 0$ for $x \in [-X, X]$, and hence there exists an $M_0 > 0$ independent of $X > -X^*$ such that $S_{w,X}(x) + E_{w,X}(x) \leq M_0$ for $x \in [-X, X]$. Therefore, $E_{w,X}(x) \leq M_0$ for $x \in [-X, X]$. Integrating (3.30) from $-X$ to X, we get

$$\int_{-X}^X E_{w,X}(x) dx \leq -D_w S'_{w,X}(-X) + (s+v_w) [S_{w,X}(-X) + E_{w,X}(-X)]$$
$$\leq -D_w p'_2(-X) + (s+v_w) [p_2(-X) + \psi_2(-X)] \leq M_0$$

for some $M_0 > 0$ independent of $X > -X^*$. In the last inequalities, we used the fact that $S'_{w,X}(X) \leq 0, E'_{w,X}(-X) \geq 0$ and $E'_{w,X}(X) \leq 0$. Consequently, by some

arguments as done for $S_{c,X}$ and $I_{c,X}$, it follows from (3.24), (3.26) and (3.27) that there exists an $M_0 > 0$, which is independent of $X > -X^*$, such that

$$\|S_{w,X}\|_{C^1[-X,X]} \le M_0, \|E_{w,X}\|_{C^1[-X,X]} \le M_0, \|I_{w,X}\|_{C^2[-X,X]} \le M_0.$$

It is easy to see that there exists an $M_0 > \sup_{x \in \mathbb{R}} \psi_5(x)$ such that M_0 is a supersolution of (3.29). Then we have $V_X(x) \le M_0$ for $x \in [-X, X]$, proving (ii).

(iii) For $Y \in (-X'', X)$, applying the L^p $(p \ge 2)$ estimates of linear elliptic differential equations to V_X and (3.29), we have

$$\|V_X\|_{W^{2,p}(U)} \le C \left(r_c \|I_{c,X}\|_{L^p(U)} + r_{ew} \|E_{w,X}\|_{L^p(U)} + r_{iw} \|I_{w,X}\|_{L^p(U)} \right.$$
$$\left. + r_d \|I_{d,X}\|_{L^p(U)} + \|\varphi\|_{W^{2,p}(U)} + \|V_X\|_{L^p(U)} \right),$$

where $U = (-Y, Y)$, $C = C(Y) > 0$ is a constant and φ is taken to be a linear function connecting the points $(-Y, V_X(-Y))$ and $(Y, V_X(Y))$. Consequently, we can find a constant $C_0 > 0$ which is only dependent of Y such that $\|V_X\|_{W^{2,p}(-Y,Y)} \le C_0$ for any $X \ge Y$. Since $W^{2,p}(-Y,Y) \hookrightarrow C^{1,\alpha}[-Y,Y]$ for $\alpha = 1 - \frac{1}{p}$, the embedding theorem further implies that there exists a constant $C > 0$ only dependent on Y such that $\|V_X\|_{C^{1,\alpha}[-Y,Y]} \le C \|V_X\|_{W^{2,p}(-Y,Y)}$. Therefore, $\|V_X\|_{C^{1,\alpha}[-Y,Y]} \le C_0$ for some $C_0 > 0$, which is only dependent on Y. From (3.29), we further have $\|V_X\|_{C^2[-Y,Y]} \le C_0$ for some $C_0 = C_0(Y) > 0$. By a similar argument, we have $\|I_{d,X}\|_{C^2[-Y,Y]} \le C_0$ for some $C_0 = C_0(Y) > 0$.

Finally, differentiating two sides of (3.23)-(3.29) gives (3.22) for some $C_0(Y)$. $\square$

We now establish our main results. Let $\{X_n\}$ be an increasing sequence with $X_n > X^*$ and $\lim_{n \to \infty} X_n = +\infty$. Then the solutions

$$(S_{c,X_n}, S_{w,X_n}, I_{c,X_n}, E_{w,X_n}, I_{w,X_n}, I_{d,X_n}, V_{X_n}) \in \Gamma_{X_n}$$

satisfy Theorem 3.6 as well as (3.23)-(3.29). We can assume (if necessary, taking a subsequence) they converge to some functions $(S_{c,*}, S_{w,*}, I_{c,*}, E_{w,*}, I_{w,*}, I_{d,*}, V_*)$ as $n \to \infty$ in the following topologies:

$$I_{c,X_n} \to I_{c,*}, I_{w,X_n} \to I_{w,*} \text{ in } C^1_{loc}(\mathbb{R}),$$
$$S_{c,X_n} \to S_{c,*}, S_{w,X_n} \to S_{w,*}, E_{w,X_n} \to E_{w,*}, I_{d,X_n} \to I_{d,*}, V_{X_n} \to V_* \text{ in } C^2_{loc}(\mathbb{R}).$$

Furthermore, $(S_{c,*}, S_{w,*}, I_{c,*}, E_{w,*}, I_{w,*}, I_{d,*}, V_*)$ satisfy system (2.1) and

$$p_1(x) \le S_{c,*}(x) \le S_{c0}, p_2(x) \le S_{w,*}(x) \le S_{w0},$$
$$\psi_1(x) \le I_{c,*}(x) \le \min\{M_0, \phi_1(x)\},$$
$$\psi_2(x) \le E_{w,*}(x) \le \min\{M_0, \phi_2(x)\}, \psi_3(x) \le I_{w,*}(x) \le \min\{M_0, \phi_3(x)\},$$
$$\psi_4(x) \le I_{d,*}(x) \le \min\{N_d, \phi_4(x)\}, \psi_5(x) \le V_*(x) \le \min\{M_0, \phi_5(x)\}$$

for $x \in \mathbb{R}$. In addition, since $S'_{c,X_n}(x) \le 0$ and $S'_{w,X_n}(x) + E'_{w,X_n}(x) \le 0$ for $x \in [-X_n, X_n]$, we have $S'_{c,*}(x) \le 0$ and $S'_{w,*}(x) + E'_{w,*}(x) \le 0$ for $x \in \mathbb{R}$.

Let $S_{c,*}(+\infty) = S_c^\infty$ and $S_{w,*}(+\infty) + E_{w,*}(+\infty) = S_w^\infty$. Then there must be $S_c^\infty < S_{c0}$ and $S_w^\infty < S_{w0}$. In fact, if $S_c^\infty = S_{c0}$, it follows that $S_{c,*}(x) = S_{c0}$ for all $x \in \mathbb{R}$, and hence $I_{c,*}(x) \equiv 0$, a contradiction. Since $\|S_{c,*}\|_{C^2(\mathbb{R})} < +\infty$, the fluctuation lemma implies that $S'_{c,*}(\pm\infty) = 0$. Hence, we have

$$\int_{-\infty}^{+\infty} [\beta_c S_{c,*}(x) V_*(x) + \alpha_c S_{c,*}(x) I_{c,*}(x)/N_{c,*}(x)] \, dx = (s + v_c)(S_{c0} - S_c^\infty) > 0.$$

It follows that $\int_{-\infty}^{\infty} I_{c,*}(x)\,dx < \infty$. By virtue of $\|I_{c,*}\|_{C^1(\mathbb{R})} < +\infty$, we have $I_{c,*}(\pm\infty) = 0$, and it turns out that

$$(3.31) \quad \int_{-\infty}^{\infty} I_{c,*}(x)\,dx = \frac{1}{d_{ic}}(s+v_c)(S_{c0} - S_c^{\infty}) \text{ and } I_{c,*}(x) \le S_{c0} - S_c^{\infty} \ \forall x \in \mathbb{R}.$$

Similarly, we can prove that

$$(3.32) \qquad \int_{-\infty}^{\infty} E_{w,*}(x)\,dx = \frac{1}{\mu_w}(s+v_w)(S_{w0} - S_w^{\infty}), \ E_{w,*}(\pm\infty) = 0,$$

$$(3.33) \qquad\qquad\qquad \int_{-\infty}^{\infty} I_{w,*}(x)\,dx = \frac{s+v_w}{d_{iw}}(S_{w0} - S_w^{\infty}),$$

$$(3.34) \qquad\qquad I_{w,*}(\pm\infty) = 0, \ I_{w,*}(x) \le \frac{s+v_w}{s}(S_{w0} - S_w^{\infty}) \ \forall x \in \mathbb{R}.$$

In conclusion, we have established the following

Theorem 3.7. *Assume that $\mathcal{R}_0 = \rho\left(\mathcal{F}\mathcal{V}^{-1}\right) > 1$. Then for any $s > s^*$, where s^* is determined by Lemma* 3.1*, system* (1.1) *admits a nonnegative traveling wave solution* $(S_{c,*}(\xi), S_{w,*}(\xi), I_{c,*}(\xi), E_{w,*}(\xi), I_{w,*}(\xi), I_{d,*}(\xi), V_*(\xi))$ *with* $\xi = x + st$ *satisfying* $S'_{c,*}(\xi) \le 0$, $S_{c,*}(-\infty) = S_{c0}$, $S_{c,*}(+\infty) = S_c^{\infty} < S_{c0}$, $S'_{w,*}(\xi) \le 0$, $S_{w,*}(-\infty) = S_{w0}$, $S_{w,*}(+\infty) = S_w^{\infty} < S_{w0}$, $I_{c,*}(\pm\infty) = 0$, $I_{d,*}(-\infty) = 0$, $\sup_{x\in\mathbb{R}} I_{d,*}(x) < N_d$, $V_*(-\infty) = 0$, $\sup_{x\in\mathbb{R}} V_*(x) < +\infty$ *and* (3.31)-(3.34).

Note that if $(S_c^{\infty})^2 + (S_w^{\infty})^2 > 0$, then by the third and fourth equations of (2.1) we have $V_*(+\infty) = 0$, which yields $V'_*(\pm\infty) = 0$ and $V''_*(\pm\infty) = 0$. Hence, by the last equation of (2.1) we have $I_{d,*}(+\infty) = 0$. Unfortunately, we cannot rule out the case where $S_c^{\infty} = S_w^{\infty} = 0$. In such a case, the traveling wave system can be reduced to the following:

$$\begin{cases} (s+v_d)\,I'_d = \beta_d\,(N_d - I_d)\,V + \alpha_d\,(N_d - I_d)\,I_d/\,N_d - \gamma_d I_d + D_d I''_d, \\ (s+v_v)\,V' = r_d I_d - (d_v + d_n)\,V + D_v V''. \end{cases}$$

Such a system can admit a positive equilibrium

$$I_d = \frac{\frac{\beta_d r_d N_d}{d_v + d_n} + \alpha_d - \gamma_d}{\frac{\beta_d r_d N_d}{d_v + d_n} + \frac{\alpha_d}{N_d}}, \quad V = \frac{r_d}{d_v + d_n} I_d.$$

Therefore, it is possible to have $I_{d,*}(+\infty) > 0$ and $V_*(+\infty) > 0$. It remains an interesting problem for future studies when $I_{d,*}(+\infty) = V_*(+\infty) = 0$.

We also remark that s^* is obtained from the stability analysis of the traveling system in the region $\chi \to -\infty$ and hence is expected to be the minimal wave speed. As the system involves multiple bird species and a component for environmental contamination (acting like a vector), detailed analysis about the minimal wave speed, its connection with the propagation speed, and the stability of wave solutions is not available at this stage.

References

1. S. Agmon, A. Douglis, L. Nirenberg, *Estimates near the boundary for solutions of elliptic differential equations satisfying general boundary conditions. I*, Comm. Pure Appl. Math. **12** (1959), 623–727. MR0125307 (23:A2610)
2. S. Agmon, A. Douglis, L. Nirenberg, *Estimates near the boundary for solutions of elliptic differential equations satisfying general boundary conditions. II*, Comm. Pure Appl. Math. **17** (1964), 35–92. MR0162050 (28:5252)
3. A. Berman, M. Neumann, Ronald J. Stern, *Nonnegative matrices in dynamic systems*, John Wiley & Sons, New York, 1989. MR1019319 (90j:93030)

4. A. Ducrot, P. Magal, *Travelling wave solutions for an infection-age structured model with diffusion*, Proc. Royal Soc. Edinburgh (A) **139** (2009), 459–482. MR2506782 (2010f:35417)
5. A. Ducrot, P. Magal, S. Ruan, *Travelling wave solutions in multigroup age-structured epidemic models*, Arch. Rational Mech. Anal. **195** (2010), 311–331. MR2564476
6. D. Gilbarg, N.S. Trudinger, *Elliptic partial differential equations of second order*, Springer, Berlin, 2001. MR1814364 (2001k:35004)
7. M. Hirsch, H. Hanish, J.-P. Gabriel, *Differential equation model of some parasitic infections: Methods for the study of asymptotic behavior*, Comm. Pure Appl. Math. **38** (1985), 733–753. MR812345 (87f:92015)
8. Y. Hosono, B. Ilyas, *Traveling waves for a simple diffusive epidemic model*, Math. Models Methods Appl. Sci. **5** (1995), 935–966. MR1359215 (96j:35248)
9. R. Liu, V.R.S.K. Duvvuri, J. Wu, *Spread pattern formation of* H5N1-*avian influenza and its implications for control strategies*, Math. Model. Nat. Phenom. **3(7)** (2008), 161–179. MR2460260 (2009j:92059)
10. Z.C. Wang, J. Wu, *Traveling waves of a diffusive Kermack-McKendrick epidemic model with non-local delayed transmission*, Proc. R. Soc. London Ser. A **466** (2010), 237–261. MR2564926 (2011b:92072)

SCHOOL OF MATHEMATICS AND STATISTICS, LANZHOU UNIVERSITY, LANZHOU, GANSU, 730000, PEOPLE'S REPUBLIC OF CHINA
E-mail address: `wangzhch@lzu.edu.cn`

CENTRE FOR DISEASE MODELING, YORK INSTITUTE FOR HEALTH RESEARCH AND DEPARTMENT OF MATHEMATICS AND STATISTICS, YORK UNIVERSITY, TORONTO, ONTARIO, M3J 1P3, CANADA
E-mail address: `wujh@mathstat.yorku.ca`

DEPARTMENT OF MATHEMATICS AND DEPARTMENT OF ZOOLOGY AND PHYSIOLOGY, UNIVERSITY OF WYOMING, LARAMIE, WYOMING 82072
E-mail address: `Rongsong.Liu@uwyo.edu`